# <span id="page-0-0"></span>SDM5008 Advanced Control for Robotics Lecture 3: Exponential Coordinate of Rigid Body Configuration

Prof. Wei Zhang

Southern University of Science and Technology, Shenzhen, China

### <span id="page-1-0"></span>**Outline**

• [Exponential Coordinate of](#page-2-0)  $SO(3)$ 

- [Euler Angles and Euler-Like Parameterizations](#page-8-0)
- [Exponential Coordinate of SE\(3\)](#page-11-0)

### <span id="page-2-0"></span>**Outline**

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# Towards Exponential Coordinate of SO(3)

- Recall the polar coordinate system of the complex plane:
	- Every complex number  $z = x + jy = \rho e^{j\phi}$
	- Cartesian coordinate  $(x, y) \leftrightarrow$  polar coorindate  $(\rho, \phi)$
	- For some applications, polar coordinate is preferred due to its geometric meaning.

\n- \n
$$
\mathsf{OM} = \{ (t, \sin(2n\pi t)) : t \in (0, 1), n = 1, 2, 3, \ldots \}
$$
\n
\n- \n
$$
\mathsf{OM} \subseteq \mathbb{R}^2
$$
\n



Exponential Coordinate of $SO(3)$	$SO(3) = \{Re\in$ $\mathbb{R}^{25}$ .
Proposition [Exponential Coordinate $\leftrightarrow SO(3)$ ]\n	$Re^{T}R \circ I$ \n
For any unit vector $[\hat{\omega}] \in so(3)$ and any $\theta \in \mathbb{R}$ ,	$\mathbb{R} \subseteq \{ker \times \text{Mult } R\} = 1$ \n
For any $R \in SO(3)$ , there exists $\hat{\omega} \in \mathbb{R}^3$ with $  \hat{\omega}   = 1$ and $\hat{\theta} \in \mathbb{R}$ such that\n	
$R = e^{[\hat{\omega}]\theta}$	$R = e^{[\hat{\omega}]\theta}$

$$
\begin{array}{ll}\n\exp: & [\hat{\omega}]\theta \in so(3) & \stackrel{\mathcal{Q} \times \mathcal{K}}{\longleftrightarrow} & R \in SO(3) \\
\log: & R \in SO(3) & \stackrel{\mathcal{Q} \times \mathcal{K}}{\longrightarrow} & [\hat{\omega}]\theta \in so(3)\n\end{array}
$$

- The vector $\left(\hat{\omega}\theta\right)$  is called the *exponential coordinate* for  $R$
- The exponential coordinates are also called the canonical coordinates of the rotation group  $SO(3)$

# Rotation Matrix as Forward Exponential Map

• Exponential Map: By definition

$$
e^{[\omega]\theta} = I + \theta[\omega] + \frac{\theta^2}{2!}[\omega]^2 + \frac{\theta^3}{3!}[\omega]^3 + \cdots
$$

• Rodrigues' Formula: Given any unit vector  $[\hat{\omega}] \in so(3)$ , we have



### Examples of Forward Exponential Map

• Rotation matrix  $R_x(\theta)$  (corresponding to  $\hat{x}\theta$ )  $dx$   $b = (a)$   $b$  $\hat{\mathcal{N}} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix} \implies [\hat{\mathcal{N}}] = \begin{bmatrix} \hat{\mathcal{D}} & \hat{\mathcal{O}} & \hat{\mathcal{O}} \\ \hat{\mathcal{O}} & \hat{\mathcal{O}} & \hat{\mathcal{O}} \\ \hat{\mathcal{O}} & \hat{\mathcal{O}} & \hat{\mathcal{O}} \end{bmatrix} \implies \hat{\mathcal{R}}_{\mathcal{N}}(\hat{\mathcal{O}}) \approx \hat{\mathcal{R}} + (\hat{\mathcal{N}}; \hat{\mathcal{O}})$  $= e^{(\lambda)\theta}$  $\Rightarrow Q_{\alpha}(9) = 1 + \sin\theta \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{bmatrix} + (1 - \cos\theta) \begin{bmatrix} 0 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos\theta & -\sin\theta \\ 0 & \sin\theta & \cos\theta \end{bmatrix}$ • Rotation matrix corresponding to  $(\underbrace{1,0,1}_{\text{CYP}})^T$ 

$$
\widehat{\omega}\theta = \begin{bmatrix} \stackrel{\bullet}{\cdot} \\ \stackrel{\bullet}{\cdot} \end{bmatrix} , \widehat{\omega} = \frac{1}{\pi} \begin{bmatrix} \stackrel{\bullet}{\cdot} \\ \stackrel{\bullet}{\cdot} \end{bmatrix} , \widehat{\theta} = \overline{12}
$$
  

$$
\begin{bmatrix} \stackrel{\bullet}{\cdot} \\ \stackrel{\bullet}{\cdot} \end{bmatrix} \longrightarrow R = 0
$$

 $A = \begin{bmatrix} a_1 \\ a_2 \end{bmatrix}$ 

# Logarithm of Rotations

• If  $R = I$ , then  $\theta = 0$  and  $\hat{\omega}$  is undefined.

• If  $tr(R) = -1$ , then  $\theta = \pi$  and set  $\hat{\omega}$  equal to one of the following

$$
\frac{1}{\sqrt{2(1+r_{33})}}\left[\begin{array}{c}r_{13}\\r_{23}\\1+r_{33}\end{array}\right], \frac{1}{\sqrt{2(1+r_{22})}}\left[\begin{array}{c}r_{12}\\1+r_{22}\\r_{32}\end{array}\right], \frac{1}{\sqrt{2(1+r_{11})}}\left[\begin{array}{c}1+r_{11}\\r_{21}\\r_{31}\end{array}\right]
$$

• Otherwise,  $\theta = \cos^{-1}\left(\frac{1}{2}(\text{tr}(R) - 1)\right) \in [0, \pi)$  and  $[\hat{\omega}] = \frac{1}{2\sin(\theta)}(R - R^T)$  $e \times p(\cdot)$ 

<span id="page-8-0"></span>

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# Euler Angle Representation of Rotation



- A common method of specifying a rotation matrix is through three independent quantities called Euler Angles.
- $\bullet$  Euler angle representation
	- Initially, frame  $\{0\}$  coincides with frame  $\{1\}$  $\mathbf{t}$
	- Rotate  $\{1\}$  about  $\hat{z}_0$  by an angle  $\alpha$ , then rotate about  $\hat{y}_a$  axis by  $\beta$ , and then if  $\alpha$ , then lotate about  $\alpha_0$  by an angle  $\alpha$ , then lotate about  $y_a$  axis by  $\rho$ , and rotate about the  $\hat{z}_b$  axis by  $\gamma$ . This yields a net orientation  ${}^0R_1(\alpha,\beta,\gamma)$ parameterized by the ZYZ angles  $(\alpha, \beta, \gamma) \Longleftrightarrow \ ^{\circ} \mathcal{R}_{l}$  (

$$
{}^{\circ}R_1(\alpha,\beta,\gamma)=\underbrace{R_z(\alpha)}_{\sim\sim\sim\sim\sim R_z}R_y(\beta)\underbrace{R_z(\gamma)}_{\circ} \overline{R_1(\alpha,\beta,\delta)} = \underbrace{R_1(\hat{s},\lambda)}_{\circ}R_2(\hat{s},\lambda)}.
$$

### Other Euler-Like Parameterizations

- Other types of Euler angle parameterization can be devised using different ordered sets of rotation axes
- Common choices include:
	- ZYX Euler angles: also called Fick angles or yaw, pitch and roll angles
	- YZX Euler angles (Helmholtz angles)



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# Exponential Map of  $se(3)$ : From Twist to Rigid Motion

**Theorem 1 [Exponential Map of**  $se(3)$ ]: For any  $\mathcal{V} = (\omega, v)$  and  $\theta \in \mathbb{R}$ , we have  $e^{|\mathcal{V}|\theta} \in SE(3)$ • Case 1  $(\omega = 0)$ :  $e^{[\mathcal{V}]\theta} = \begin{bmatrix} I & v\theta \\ 0 & 1 \end{bmatrix}$ 

• Case 2 ( $\omega \neq 0$ ): without loss of generality assume  $\|\omega\| = 1$ . Then

$$
e^{[\mathcal{V}]\theta} = \begin{bmatrix} e^{[\omega]\theta} & G(\theta)v \\ 0 & 1 \end{bmatrix}, \text{ with } G(\theta) = I\theta + (1 - \cos(\theta))[\omega] + (\theta - \sin(\theta))[\omega]^2 \quad (1)
$$
  
For any  $t\omega^t s^t$   $\begin{bmatrix} \sqrt{2} & \theta & \sqrt{2} \\ \theta & \sqrt{2} & \theta & \sqrt{2} \\ 0 & \sqrt{2} & \sqrt{2} \end{bmatrix} \begin{bmatrix} \theta & \sqrt{2} & \theta \\ \theta & \sqrt{2} & \theta \end{bmatrix}$ 

# Log of  $SE(3)$ : from Rigid-Body Motion to Twist

**Theorem 2 [Log of**  $SE(3)$ ]: Given any  $T = (R, p) \in SE(3)$ , one can always find twist  $\mathcal{S} = (\omega, v)$  and a scalar  $\theta$  such that

$$
e^{[\mathcal{S}]\theta} = T = \left[ \begin{array}{cc} R & p \\ 0 & 1 \end{array} \right]
$$

#### Matrix Logarithm Algorithm:

- If  $R = I$ , then set  $\omega = 0$ ,  $v = p/||p||$ , and  $\theta = ||p||$ .
- Otherwise, use matrix logarithm on  $SO(3)$  to determine  $\omega$  and  $\theta$  from R. Then v is calculated as  $v=G^{-1}(\theta)p$ , where

$$
G^{-1}(\theta) = \frac{1}{\theta}I - \frac{1}{2}[\omega] + \left(\frac{1}{\theta} - \frac{1}{2}\cos\frac{\theta}{2}\right)[\omega]^2
$$

# Exponential Coordinates of Rigid Transformation

• To sum up, screw axis  $\mathcal{S} = (\omega, v)$  can be expressed as a normalized twist; its matrix representation is

$$
[\mathcal{S}] = \left[ \begin{array}{cc} [\omega] & v \\ 0 & 0 \end{array} \right] \in se(3)
$$

- A point started at  $p(0)$  at time zero, travel along screw axis S at unit speed for time  $t$  will end up at  $\tilde{p}(t) = e^{[\mathcal{S}]}{}^t \tilde{p}(0)$
- Given  $S$  we can use Theorem 1 to compute  $e^{[S]t} \in SE(3)$ ;
- Given  $T \in SE(3)$ , we can use Theorem 2 to find  $S = (\omega, v)$  and  $\theta$  such that  $e^{[\mathcal{S}]\theta}=T.$
- We call  $S\theta$  the **Exponential Coordinate** of the homogeneous transformation  $T \in SE(3)$  $T \in SE(3)$  $T = e^{1370}$  $SO \in \mathbb{R}^6$

# More Space

# <span id="page-16-0"></span>More Space